

# Photoluminescent behavior of $\text{SrB}_4\text{O}_7:\text{RE}^{2+}$ ( $\text{RE} = \text{Sm}$ and $\text{Eu}$ ) prepared by Pechini, combustion and ceramic methods

R. Stefani<sup>a</sup>, A.D. Maia<sup>a</sup>, E.E.S. Teotonio<sup>a</sup>, M.A.F. Monteiro<sup>a</sup>, M.C.F.C. Felinto<sup>b</sup>,  
H.F. Brito<sup>a,\*</sup>

<sup>a</sup>*Instituto de Química, Universidade de São Paulo, C.P. 26077, CEP 05513-970 São Paulo-SP, Brazil*

<sup>b</sup>*Instituto de Pesquisas Energéticas e Nucleares, Av. Prof. Lineu Prestes 2242, Cidade Universitária, São Paulo-SP, Brazil*

Received 28 October 2005; received in revised form 27 December 2005; accepted 3 January 2006

Available online 2 February 2006

## Abstract

This work reports the preparation of system containing  $\text{RE}^{2+}$  ions ( $\text{RE} = \text{Sm}$  and  $\text{Eu}$ )-doped in  $\text{SrB}_4\text{O}_7$  matrix by ceramic, Pechini and combustion methods. These compounds were prepared by reduction of  $\text{RE}^{3+}$  to  $\text{RE}^{2+}$  in air, which exhibit some different features according to the preparation method. Photoluminescent properties of these systems were investigated based on the emission and excitation spectral data. The emission spectra of  $\text{SrB}_4\text{O}_7:\text{Eu}^{2+}$  system prepared by combustion and Pechini methods are characterized by a broad band assigned to interconfigurational  $4f^6 5d \rightarrow 4f^7$  transition, while  $\text{SrB}_4\text{O}_7:\text{Sm}^{2+}$  compound exhibit narrow emission bands arising from intraconfigurational- $4f^6$  also shows  $^4G_{5/2} \rightarrow ^6H_J$  transitions ( $J' = \frac{5}{2}, \frac{7}{2}, \frac{9}{2}$  and  $\frac{11}{2}$ ) arising from  $\text{Sm}^{3+}$  ion, transitions.  $\text{SrB}_4\text{O}_7:\text{RE}$  system prepared by combustion method presents emission bands from  $\text{RE}^{3+}$  ions as intense as that arising from  $\text{RE}^{2+}$ , suggesting that the preparation route is not efficient for the reduction  $\text{RE}^{3+} \rightarrow \text{RE}^{2+}$  process. Emission quantum efficiency and radiative emission rates of  $\text{Sm}^{2+}$  ion are determined and discussed.

© 2006 Elsevier Inc. All rights reserved.

**Keywords:** Luminescence; Borates; Divalent rare earths; Luminescence;  $f-d$  transitions

## 1. Introduction

In recent years, the versatility of efficient luminescent materials have attracted researchers' attention due to their applications in the field of fluorescent lamps, display devices, detector systems represented by X-ray screens, immunoassays and scintillates of phosphor marking [1–4]. Borate systems have proved to be potential candidates for aforementioned applications [4,5]. For example, strontium tetraborate ( $\text{SrB}_4\text{O}_7$ ) has been investigated as a potential non-linear optical material with some excellent mechanical and optical properties such as high mechanical strength, non-hygroscopicity and high optical damage threshold [6].

$\text{Eu}^{2+}$  and  $\text{Gd}^{3+}$  ions ( $4f^7$ -configuration) present  $^8S_{7/2}$  and  $^6P_{7/2}$  levels as ground and excited states, respectively.

Consequently, a qualitative interpretation of intraconfigurational transitions can be made in a similar way, but the position of excited  $^6P_{7/2}$  level of  $\text{Eu}^{2+}$  ion is at a lower energy ( $\sim 28,000 \text{ cm}^{-1}$ ) than  $\text{Gd}^{3+}$  ions ( $\sim 32,000 \text{ cm}^{-1}$ ). This behavior is also observed for  $\text{Sm}^{2+}$  and  $\text{Eu}^{3+}$  ions that present  $4f^6$ -configuration where the first excited  $^5D_0$  level is around  $14,500$  and  $17,300 \text{ cm}^{-1}$ , respectively. Fig. 1 shows partial energy level diagram of free  $\text{Sm}^{2+}$ ,  $\text{Sm}^{3+}$ ,  $\text{Eu}^{2+}$ ,  $\text{Eu}^{3+}$  and  $\text{Gd}^{3+}$  ions including the lowest energy  $f-d$  levels [7–10].

The most important characteristic of divalent europium ion is that, it shows broad emission bands, in visible spectrum, arising from interconfigurational  $4f^7 \rightarrow 4f^6 5d^1$  transitions which are strongly dependent on chemical environment since  $5d$  orbitals have a great radial extension and are more sensitive to the ligand field [11–13]. Accordingly,  $f-d$  emission bands can be shifted from blue to red region, which makes  $\text{Eu}^{2+}$  ion potentially promising

\*Corresponding author. Fax: +55 11 3815 5579.

E-mail address: [hefbrito@iq.usp.br](mailto:hefbrito@iq.usp.br) (H.F. Brito).

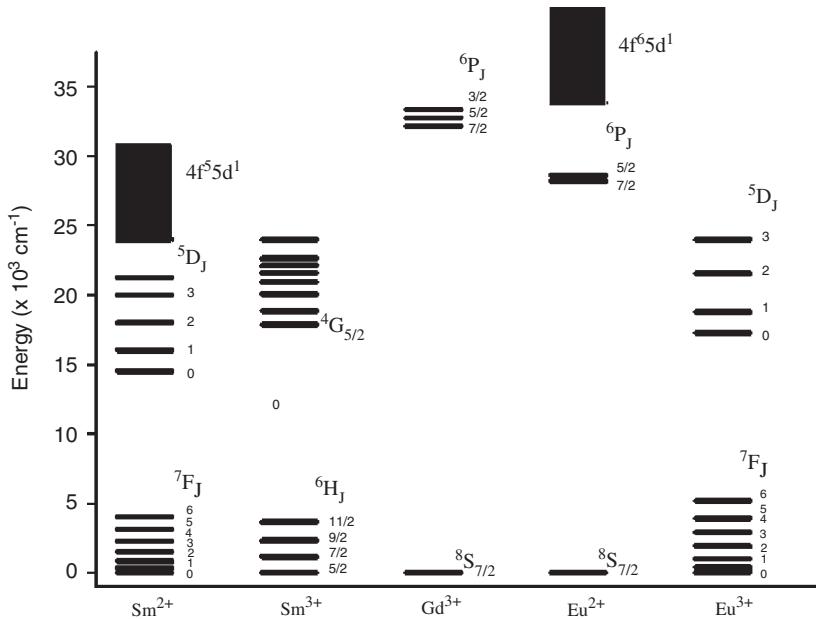


Fig. 1. Partial energy level diagram of free divalent and trivalent rare earth ions ( $\text{Sm}^{2+}$ ,  $\text{Sm}^{3+}$ ,  $\text{Eu}^{2+}$ ,  $\text{Eu}^{3+}$  and  $\text{Gd}^{3+}$ ) containing the lowest energy  $f$ - $d$  levels.

systems as phosphors for applications in displays, lamps and luminescent paintings used for several signs and directions, meaning to provide safe traffic, road markers and important devices.

On the other hand, emission spectra of divalent samarium ions, in the spectral range from 770 to 800 nm, are generally dominated by narrow bands that are assigned to the intraconfigurational ( $4f^6$ – $4f^6$ ) transitions. The presence of these sharp bands is due to  $4f$  electrons that are well protected from ligand field by  $5s^25p^6$  external orbitals and band positions are only slightly affected by environment around samarium ion [13].

Solid state systems containing divalent rare earth ions have been prepared using reducing agents such as  $\text{H}_2$  to reduce  $RE^{3+}$  to  $RE^{2+}$  ions. However, some  $RE^{2+}$ -systems can also be prepared by reduction process of  $RE^{3+}$  ion using solid state reaction at high temperature in air and/or  $\text{N}_2$  atmospheres. For example, the  $\text{Eu}^{3+}$  or  $\text{Sm}^{3+}$  ions doped into  $\text{SrB}_4\text{O}_7$  were reduced to their corresponding divalent rare earth ions by a solid state reaction at high temperature in air [14]. The ease of incorporation and stabilization of  $\text{Eu}^{2+}$  and  $\text{Sm}^{2+}$  ions in this material is due to unique structure of  $\text{SrB}_4\text{O}_7$  matrix considering that all boron atoms are tetrahedrally coordinated forming a three-dimensional  $(\text{B}_4\text{O}_7)_{\infty}$ -network by corner-sharing, containing channels parallel to  $b$ -axis. The strontium ion fits into these channels and is surrounded by nine oxygen atoms, giving  $\text{SrO}_9$  polyhedral, with  $C_s$  symmetry. Divalent rare earth ions are therefore located in a “cage” formed by  $\text{BO}_4$  unit from  $(\text{B}_4\text{O}_7)_{\infty}$ -network with rigid structure which prevents oxidation of  $RE^{2+}$  to  $RE^{3+}$  ions [15].

The purpose of this work is to investigate photoluminescence properties of  $\text{Eu}^{2+}$  and  $\text{Sm}^{2+}$  doped in  $\text{SrB}_4\text{O}_7$  host

prepared by Pechini [16] and combustion [17] as alternative methods, in milder experimental conditions than the solid state reaction (ceramic method) [18]. Luminescent behavior of  $\text{Sm}^{2+}$  and  $\text{Eu}^{2+}$  ions in this matrix were investigated based on intra and interconfigurational transitions. Experimental quantum efficiencies ( $\eta$ ) for the emitting  $^5D_0$  level of  $\text{Sm}^{2+}$  ion were determined and discussed.

## 2. Experimental

The preparation of  $\text{SrB}_4\text{O}_7:RE^{2+}$  system ( $RE = \text{Sm}$  and  $\text{Eu}$ ) by ceramic method was carried out in air. Stoichiometric amounts of  $\text{Sr}(\text{NO}_3)_2$  (Synth) and  $\text{H}_3\text{BO}_3$  (Cinética Química) together with a 3 mol% excess of  $\text{H}_3\text{BO}_3$  to compensate for  $\text{B}_2\text{O}_3$  evaporation were thoroughly mixed and grounded. The dopant concentration of  $RE^{3+}$  added as  $RE(\text{NO}_3)_3 \cdot 6\text{H}_2\text{O}$  (prepared from  $\text{Eu}_2\text{O}_3$ , Aldrich, 99.9%), was 1 mol% of  $\text{Sr}^{2+}$  ion concentration in host compound. The mixtures were heated in air at  $700^{\circ}\text{C}$  in alumina crucibles for 5 h and then grounded and heated at  $850^{\circ}\text{C}$  for another 5 h. The second heating step was repeated once.

In preparation by Pechini method, the stoichiometric quantities of  $\text{Sr}(\text{NO}_3)_2$ ,  $RE(\text{NO}_3)_3 \cdot 6\text{H}_2\text{O}$  (1 mol% of  $\text{Sr}^{2+}$ ) and  $\text{H}_3\text{BO}_3$  precursors were dissolved in deionized water. In this solution, preheated at  $70^{\circ}\text{C}$ , citric acid was added in molar ratio of 3:1 (citric acid: $\text{Sr}^{2+}$ ) and ethylene glycol in weight ratio of 3.6:1 (citric acid:ethylene glycol) under stirring. Then, this mixture was heated at  $110^{\circ}\text{C}$ , and the formation of brown transparent glassy polymeric resin was observed. After, the resin was heated at  $300^{\circ}\text{C}$  for 2 h, resulting in a fluffy black mass, which was grounded into a powder. Finally, this material was calcinated at different

temperatures and times and the optimized condition was obtained by heating at 900 °C during 3 h, which produced a  $\text{SrB}_4\text{O}_7:\text{Sm}^{2+}$  powder sample.

For the system prepared by combustion method, similar molar quantities of reagents to traditional ceramic method were used. Urea in 2.5 mol% of  $\text{Sr}^{2+}$  was used as combustible, which was added to mixture of rare earth and strontium nitrates and boric acid dissolved in deionized water. The resultant solution was heated at 300 °C (3 °C/min) in the hot plate to dehydrate and decompose. Then, in a few seconds, the mixture ignites, leading to smooth deflagration with enormous swelling, producing white foam. In this chemical reaction,  $\text{NO}_2$ ,  $\text{H}_2\text{O}$ ,  $\text{NH}_3$ ,  $\text{CO}_2$ ,  $\text{O}_2$  and  $\text{N}_2$  gases are released. The voluminous and foamy solid was milled and fired slowly at 400 °C for 1 h, then 500 °C for 2 h, yielding the precursor powder of  $\text{SrB}_4\text{O}_7:\text{RE}^{2+}$ . Thermal treatment was performed at different temperatures, but the final powder was only obtained after heating at least at 700 °C for 2 h.

The purity and compositions of the samples were checked by X-ray diffraction (XRD) patterns recorded on a Philips diffractometer model X'Pert-MPD using  $\text{CuK}\alpha$  radiation (40 kV and 40 mA) in the interval of 2–70° (2θ) and 1 s of pass time, using the powder XRD method.

Photoluminescence measurements were performed in a spectrofluorimeter (SPEX-Fluorolog 2) with double grating 0.22 m monochromators (SPEX 1680), and a 450 W Xenon lamp as excitation source. A spectroscopy computer DM3000F controlled this apparatus. The excitation and emission spectra at room and liquid nitrogen temperatures were collected at an angle of 22.5° (front face).

The infrared absorption spectra were recorded in KBr pellets on a Bomen Model MB-100 spectrophotometer in the range of 4000–400  $\text{cm}^{-1}$ .

The media crystallite size of the samples was determined using the Scherrer's formula [19].

### 3. Results and discussion

IR spectra of the samples are presented in Fig. 2. The spectra of the systems prepared by ceramic and Pechini methods exhibit very similar profiles, showing that the exact positions of all the prominent bands in the host borates spectra are coincident. It is also important to report that these spectra present characteristic bands centered at 1100  $\text{cm}^{-1}$ , which are associated to tetrahedrally coordinated B atoms [20]. This result indicates that the same type of system was obtained in spite of the preparation method. On the other hand, the samples prepared by combustion method exhibit IR spectra with additional bands (above 1100  $\text{cm}^{-1}$ ), which are assigned to  $\text{BO}_3$ -trigonal group in metaborate phase. These results are confirmed by the powder XRD patterns of  $\text{SrB}_4\text{O}_7:\text{RE}^{2+}$  shown in Fig. 3. For example, it is observed that, when ceramic and Pechini methods are used, most of intense reflections are assigned to the tetraborate phase [21], while a very low proportion of metaborate is observed when the samples are prepared by the combustion method where metaborate phase is the foremost one. It is important to report that no significant difference was observed in the IR and powder XRD patterns for  $\text{SrB}_4\text{O}_7:\text{RE}^{2+}$  compound obtained by combustion method under heating at 700, 850 and 900 °C. This result gives good evidence that the same product was obtained at different temperatures.

Average crystallite size of samples prepared by different methods was determined using Scherrer's formula [19]. Values of particle sizes of divalent samarium and europium ions are 81 and 79 nm, prepared by ceramic method, 77 and 67 nm by Pechini method and 38 and 39 nm by combustion method, respectively. These results show that average particle sizes of  $\text{SrB}_4\text{O}_7:\text{RE}^{2+}$  system prepared by combustion method are almost twice smaller than those synthesized by ceramic and Pechini methods.

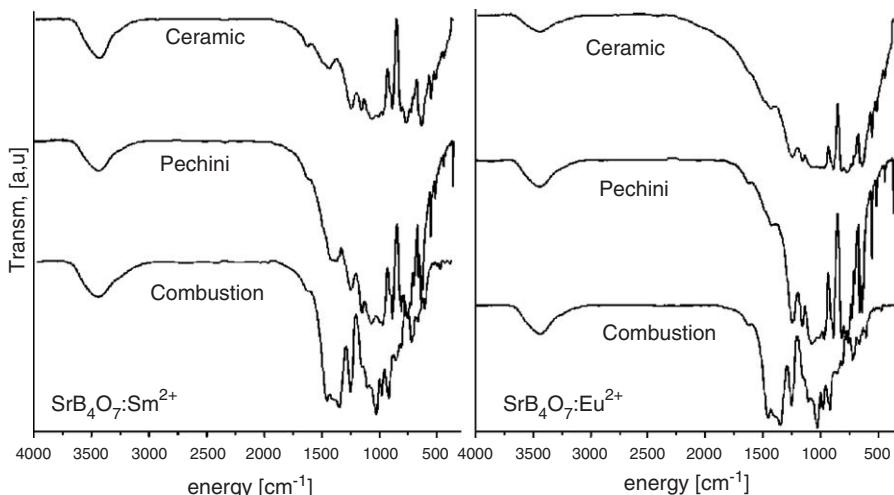


Fig. 2. Infrared spectra of  $\text{SrB}_4\text{O}_7:\text{RE}^{2+}$  (Sm and Eu) systems prepared by ceramic, Pechini and combustion methods.

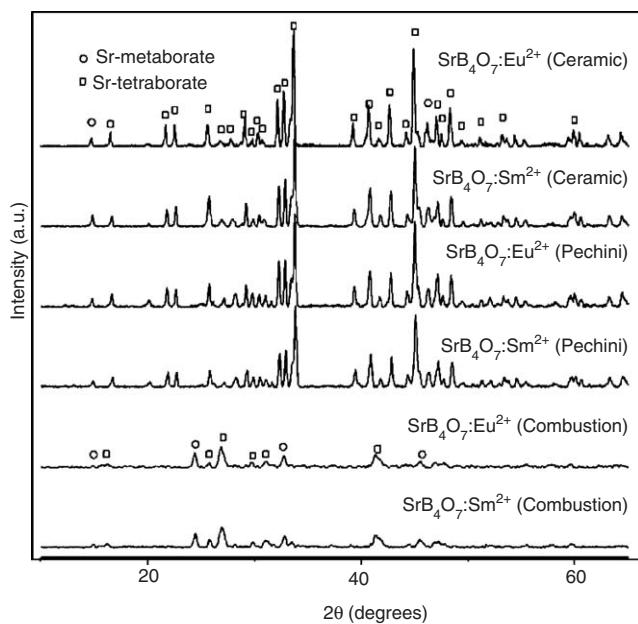


Fig. 3. X-ray diffraction patterns of  $\text{SrB}_4\text{O}_7:\text{RE}^{2+}$  (Sm and Eu) systems prepared by Pechini, combustion and ceramic methods at 900, 700 and 850 °C, respectively.

### 3.1. Photoluminescence investigation

Excitation spectra of  $\text{SrB}_4\text{O}_7:\text{Sm}^{2+}$  samples prepared by ceramic, Pechini and combustion methods recorded at 77 K are shown in Fig. 4a–c. These spectra were recorded monitoring the emission intensity of intraconfigurational  $^5D_0 \rightarrow ^7F_0$  transition (around 685 nm) centered on  $\text{Sm}^{2+}$  ion. Generally, spectra of  $\text{Sm}^{2+}$  systems display broad bands in spectral range of 250–550 nm corresponding to interconfigurational  $4f^6 \rightarrow 4f^55d$  transition, presenting two split components  $t_{2g}$  and  $e_g$  of 5d excited state from rare earth ion. A number of sharp peaks due to crystal-field interaction of  $4f^6$  are also observed. When a comparison is made among excitation spectra of  $\text{SrB}_4\text{O}_7:\text{Sm}^{2+}$  samples prepared by different methods, it is observed that samples prepared by ceramic and Pechini methods show similar excitation spectral profiles (Figs. 4a and b). On the other hand, excitation spectra of  $\text{Sm}^{2+}$ -samples prepared by combustion method exhibit an intense band around 280 nm which may be assigned to transitions of  $\text{Sm}^{3+}$  ion (Fig. 4c), since this ion was not totally converted to  $\text{Sm}^{2+}$  ion as a consequence of the presence of metaborate phase.

Figs. 4d–f exhibit emission spectra of Sm-samples recorded at 77 K with excitation monitored at 367 nm. These spectra present narrow emission bands in the range from 680 to 775 nm, that are attributed to intraconfigurational  $^5D_0 \rightarrow ^7F_J$  transitions ( $J = 0, 1, 2$  and  $3$ ) arising from  $\text{Sm}^{2+}$  ion (Table 1). These transitions are split in a maximum of  $(2J + 1)$ -components suggesting the presence of only one site of symmetry around the divalent samarium ion. Furthermore, the emission spectrum of Sm-sample prepared by combustion method also shows  $^4G_{5/2} \rightarrow ^6H_J$

transitions ( $J' = \frac{5}{2}, \frac{7}{2}, \frac{9}{2}$  and  $\frac{11}{2}$ ) arising from  $\text{Sm}^{3+}$  ion, which indicates only a partial conversion to  $\text{Sm}^{2+}$  ion and suggests a very low proportion of tetraborate phase. Since the excitation is monitored on interconfigurational  $4f^6 \rightarrow 4f^55d$  transition, the absence of broad band from emission spectra is an indicative of an efficient energy transfer from the states of  $4f^55d$  configuration to excited levels of  $\text{Sm}^{2+}$  ion.

It is important to report that the emission spectra of all Sm-samples also exhibit narrow bands arising from  $^5D_1 \rightarrow ^7F_J$  transitions in spectral range 620–675 nm. However, their intensities are higher when spectra were recorded at 77 K than at room temperature. This result is a consequence of the small energy gap between  $^5D_0$  and  $^5D_1$  excited states ( $\Delta E \approx 1350 \text{ cm}^{-1}$ ), which is very close to vibration mode of  $\text{BO}_4$  around  $1300 \text{ cm}^{-1}$ . Thus emission bands from  $^5D_1$  excited level may be explained by a thermal population of  $^5D_1$  level from  $^5D_0$  state.

Luminescence decay curves of  $\text{SrB}_4\text{O}_7:\text{Sm}^{2+}$  system were recorded monitoring emission on  $^5D_0 \rightarrow ^7F_0$  transition ( $\sim 685 \text{ nm}$ ) and excitation at 367 nm, at room temperature. The decay curves for all samples present mono-exponential behavior with similar values of lifetimes for the emitting  $^5D_0$  level (Table 2). These data are in accordance with the presence of only one non-radiative deactivation channel acting on the emitting level  $^5D_0$ .

Photoluminescence properties of  $\text{SrB}_4\text{O}_7:\text{RE}$  ( $\text{RE} = \text{Sm}$  and Eu) systems were investigated based on excitation and emission spectra recorded at room and at liquid nitrogen temperatures. All spectra recorded at low temperature are similar profiles to those obtained at room temperature, however presenting better resolution.

In order to investigate spectroscopic properties of  $\text{Sm}^{2+}$  doped in  $\text{SrB}_4\text{O}_7$  host, the transition probabilities for zero-phonon lines and emission quantum efficiency were determined. The spontaneous emission coefficients  $A_{0J}$  for  $^5D_0 \rightarrow ^7F_J$  transitions were determined in a similar procedure used for  $\text{Eu}^{3+}$  ion [20,21], by taking magnetic dipole  $^5D_0 \rightarrow ^7F_1$  transition as standard, since the emission probability of this transition is practically insensitive to chemical environment around  $\text{Eu}^{3+}$  and  $\text{Sm}^{2+}$  ions ( $\text{Eu}^{3+}$ ,  $A_{01} = 50 \text{ s}^{-1}$  and  $\text{Sm}^{2+}$ ,  $A_{01} = 29 \text{ s}^{-1}$ ). The emission coefficients were determined by following equation [22]:

$$A_{0\lambda} = A_{01} \left( \frac{S_{0\lambda}}{S_{01}} \right) \left( \frac{\sigma_{01}}{\sigma_{0\lambda}} \right), \quad (1)$$

where  $S_{0\lambda}$  and  $\sigma_{0\lambda}$  are the area under the emission band and the energy barycenter of  $^5D_0 \rightarrow ^7F_\lambda$  transition, respectively. By the relation between the lifetime of  $^5D_0$  emitting state and total decay rate ( $1/\tau = A_{\text{rad}} + A_{\text{nrad}}$ ), where  $A_{\text{rad}}$  and  $A_{\text{nrad}}$  are the radiative and non-radiative rates, respectively, the value can be calculated by following equation [23]:

$$\eta = \frac{A_{\text{rad}}}{A_{\text{rad}} + A_{\text{nrad}}}. \quad (2)$$

Table 2 presents the experimental intensity parameter for  $\text{SrB}_4\text{O}_7:\text{Sm}$  prepared by the three different

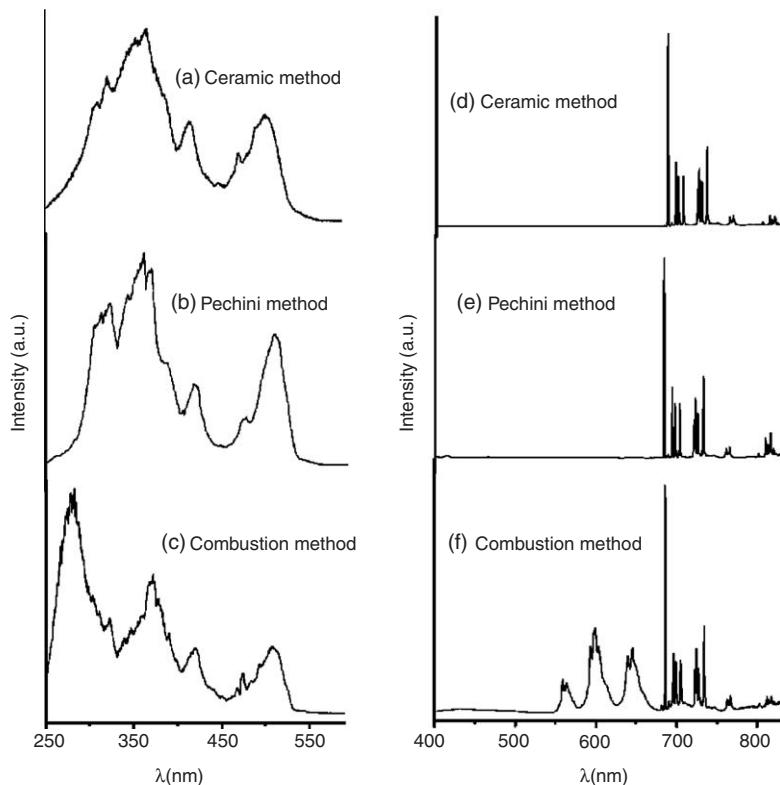


Fig. 4. Luminescence spectra of  $\text{SrB}_4\text{O}_7:\text{Sm}^{2+}$  system prepared by ceramic, Pechini and combustion methods: (a) excitation spectra recorded with emission monitored around 685 nm and (b) emission spectra recorded under excitation at 350 nm.

Table 1  
Transition energies of  $^5D_0 \rightarrow ^7F_J$  manifolds (in  $\text{cm}^{-1}$ ) observed in the emission spectra of  $\text{SrB}_4\text{O}_7:\text{Sm}^{2+}$  samples, at 77 K

Transitions	Preparation methods		
	Ceramic	Pechini	Combustion
$^5D_0 \rightarrow ^7F_0$	14,585 14,376	14,587 14,378	14,590 14,380
$^5D_0 \rightarrow ^7F_1$	14,310 14,188 13,846	14,306 14,194 13,840	14,310 14,192 13,846
$^5D_0 \rightarrow ^7F_2$	13,808 13,751 13,627	13,812 13,755 13,623	13,808 13,751 13,627
$^5D_0 \rightarrow ^7F_3$	13,140 13,044 12,465 12,330	13,131 13,046 12,493 12,358	13,113 13,041 12,456 12,315
$^5D_0 \rightarrow ^7F_4$	12,291 12,240 12,186	12,318 12,264 12,216	12,282 12,228 12,183

values has been observed when the particles sizes were changed. An important feature to be observed is the higher values of  $R_{02}$  parameter, which is given by the integral ratio between the intensities. The values of  $R_{02}$  for  $\text{Sm}^{2+}$ -systems around 0.7 are higher than those obtained for  $\text{Eu}^{3+}$  compounds, suggesting that there is a higher  $J$ -mixing in the  $\text{Sm}^{2+}$ -systems.

The luminescence properties of  $\text{SrB}_4\text{O}_7:\text{Eu}^{2+}$  compound were investigated based on excitation and emission spectral data. Figs. 5a–c show excitation spectra of europium samples recorded at 77 K by monitoring emission of  $\text{Eu}^{2+}$  ion ( $4f^65d \rightarrow 4f^7$  transition) at 367 nm. It is clearly stated that these spectra display mainly two broad bands with maxima around 270 and 310 nm that correspond to splitting of  $5d$  orbital in  $t_{2g}$  and  $e_g$  components.

Figs. 5d–f show emission spectra of  $\text{Eu}^{3+}$ -doped  $\text{SrB}_4\text{O}_7$  under excitation at 265 nm. The main features of the spectra of compounds prepared by Pechini and ceramic methods (Figs. 5d and e) present broad bands around 367 nm, which is assigned to interconfigurational transition  $4f^7 \rightarrow 4f^65d$  arising from  $\text{Eu}^{2+}$  ion. It is also observed very weak emission bands in 570–720 nm spectral range, corresponding to intraconfigurational  $^5D_0 \rightarrow ^7F_J$  transitions ( $J = 0, 1, 2, 3, 4$ ) centered on  $\text{Eu}^{3+}$  ion [23]. This result is a good evidence of the efficient reduction  $\text{Eu}^{3+} \rightarrow \text{Eu}^{2+}$  process. On the other hand, the sample prepared by combustion method (Fig. 5f) exhibits emission spectrum with transition  $^5D_0 \rightarrow ^7F_J$  ( $\text{Eu}^{3+}$ ) as intense as that arising

preparation methods. As it can be seen, quantum emission efficiency of emitting  $^5D_0$  level is around 37% suggesting that these compounds present great potential for application as phosphors. Besides, no significant change in the  $\eta$

Table 2

Emission quantum efficient ( $\eta$ ), lifetimes ( $\tau$ ) and radiative decay rates ( $A_{\text{rad}}$ ), non-radiative ( $A_{\text{nrad}}$ ) and total ( $A_{\text{total}}$ ) for  $\text{SrB}_4\text{O}_7:\text{Sm}^{2+}$  system prepared by ceramic, Pechini and combustion methods

$\text{SrB}_4\text{O}_7:\text{Sm}^{2+}$	$A_{\text{rad}} (\text{s}^{-1})$	$A_{\text{nrad}} (\text{s}^{-1})$	$A_{\text{total}} (\text{s}^{-1})$	$R_{0-2}$	$\tau (\text{ms})$	$\eta (\%)$
Ceramic	100	175	275	0.736	3.627	36
Pechini	118	194	312	0.690	3.196	37
Combustion	120	202	322	0.524	3.100	37

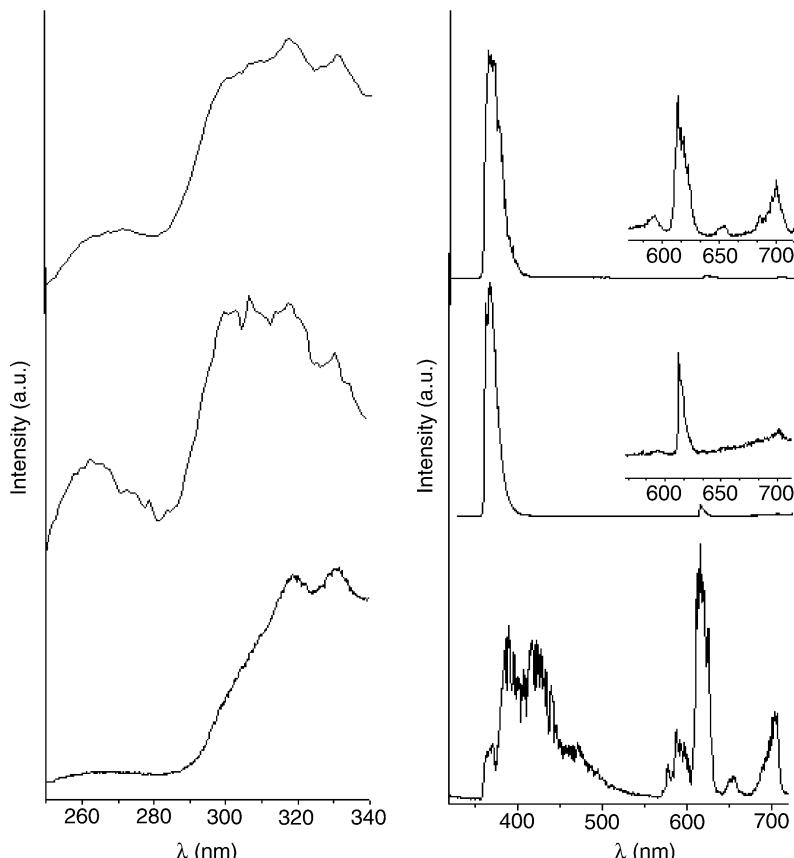


Fig. 5. Luminescence spectra of  $\text{SrB}_4\text{O}_7:\text{Eu}^{2+}$  system prepared by ceramic, Pechini and combustion methods: (a) excitation spectra recorded with emission monitored around 367 nm and (b) emission spectra recorded with excitation monitored around 265 nm.

from  $4f^65d \rightarrow 4f^7$  transition ( $\text{Eu}^{2+}$ ), indicating that the reduction process is incomplete. Similarly to samarium systems this behavior may be explained as a consequence of the structural change, suggesting a very low proportion of tetraborate phase.

#### 4. Conclusions

$\text{SrB}_4\text{O}_7:\text{RE}$  system prepared by ceramic and Pechini methods exhibits emission bands arising from  $\text{Sm}^{2+}$  and  $\text{Eu}^{2+}$  ions, suggesting an efficient reduction process from  $\text{RE}^{3+}$  to  $\text{RE}^{2+}$ . While the  $\text{SrB}_4\text{O}_7:\text{RE}$  system prepared by combustion method exhibits emission bands from  $\text{RE}^{3+}$  ions suggesting that the reduction  $\text{RE}^{3+} \rightarrow \text{RE}^{2+}$  process was incomplete.

Sm-samples present narrow emission bands attributed to intraconfigurational  $5D_0 \rightarrow 7F_J$  transitions of  $\text{Sm}^{2+}$  ion

except for the sample prepared by combustion method that also shows  $4G_{5/2} \rightarrow 6H_J$  transitions of  $\text{Sm}^{3+}$  ion. In case of  $\text{Eu}^{3+}$ -doped  $\text{SrB}_4\text{O}_7$  prepared by Pechini and ceramic methods present emission broad bands assigned to interconfigurational transition  $4f^65d \rightarrow 4f^7$  arising from  $\text{Eu}^{2+}$  ion when compared with combustion method. Based on optical data of  $\text{SrB}_4\text{O}_7:\text{RE}^{2+}$ , it is observed that the reduction of  $\text{RE}^{3+} \rightarrow \text{RE}^{2+}$  is more efficient for the system prepared by ceramic and Pechini methods than that by combustion. This result may be explained as a consequence of preferential metaborate phase for the  $\text{SrB}_4\text{O}_7:\text{RE}^{2+}$  systems when they were obtained by combustion method.

#### Acknowledgments

The authors thank the Brazilian agencies FAPESP, CNPq and CNPq-RENAMI for the financial support.

## References

- [1] B. Lei, Y. Liu, J. Liu, Z. Ye, C. Shi, *J. Solid State Chem.* 177 (2004) 1333.
- [2] T. Matsuzawa, Y. Aoki, N. Takeuchi, Y. Murayama, *J. Electrochem. Soc.* 143 (1996) 2670.
- [3] W.O. Gordon, J.A. Carter, B.M. Tissue, *J. Lumin.* 108 (2004) 339.
- [4] G. Bertrand-Chadeyron, M. El-Ghazzi, D. Boyer, R. Mahiou, J.C. Cousseins, *J. Alloys Compd.* 317 (2001) 183.
- [5] A. Diaz, D.A. Keszler, *Mater. Res. Bull.* 31 (1996) 147.
- [6] Y.U. Oseledchik, A.L. Prosvirnin, V.V. Starshenko, V.V. Osadchuk, A.I. Pisarevsky, S.P. Belokrys, A.S. Korol, N.V. Svitanko, A.F. Selevich, S.A. Krikunov, *J. Cryst. Growth* 135 (1994) 373.
- [7] P. Dorenbos, *J. Lumin.* 104 (2003) 239.
- [8] W.C. Martin, R. Zalubas, L. Hagan, *Atomic Energy levels—The Rare-Earth Elements*, National Bureau of Standards, US government Printing Office, Washington, 1978.
- [9] H.F. Brito, G.K. Liu, *J. Chem. Phys.* 112 (2000) 4334.
- [10] W.T. Carnall, G.L. Goodman, K. Rajnak, R.S. Rana, *J. Chem. Phys.* 90 (1989) 3443.
- [11] J.M. Rey, J.R.G. Thorne, R.G. Denning, H. Bill, *J. Phys.: Condens. Matter* 17 (2005) 1861.
- [12] M.F. Reid, L. van Pieterson, A. Meijerink, *J. Alloys Compd.* 344 (2002) 240.
- [13] L. Ning, C. Duan, S. Xia, M.F. Reid, P.A. Tanner, *J. Alloys Compd.* 366 (2004) 34.
- [14] Z. Pei, Q. Su, J. Zhang, *J. Alloys Compd.* 198 (1993) 51.
- [15] K. Machida, G. Adachi, J. Shiokawa, *Acta Crystallogr. B* 36 (1980) 2008.
- [16] M.U. Pechini, US Patent No. 3330697, 1967.
- [17] S. Shikao, W. Jiye, *J. Alloys Compd.* 327 (2001) 82.
- [18] G.Y. Adachi, K.I. Machida, J. Shiokawa, *J. Less-Common Met.* 93 (1983) 389.
- [19] B.D. Cullity, *Elements of X-ray Diffraction*, Addison-Wesley Publishing Company Inc., Reading, MA, 1967.
- [20] P.B. Hart, S.E.F. Smallwood, *J. Inorg. Nucl. Chem.* 24 (1962) 1047.
- [21] A. Lavat, C. Graselli, M. Santiago, J. Pomarico, E. Caselli, *Cryst. Res. Technol.* 39 (2004) 840.
- [22] C.A. Kodaira, H.F. Brito, O.L. Malta, O.A. Serra, *J. Lumin.* 101 (2003) 11.
- [23] E.E.S. Teotonio, M.C.F.C. Felinto, H.F. Brito, O.L. Malta, A.C. Trindade, R. Najjar, W. Strek, *Inorg. Chim. Acta* 357 (2004) 451.